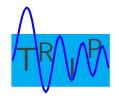
Solutions to Integro-Differential Evolution Equations with Discontinuous Coefficients

Kirk Blazek

The Rice Inversion Project Department of Computational and Applied Mathematics Rice University

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- Current: TRIP-VIGRE Postdoc in CAAM dept. at Rice University Advisor: W. W. Symes
- PhD in mathematics from the University of Washington in 2006
 Advisor: K. P. Bube The One-Dimensional Seismic Inverse Problem on a Viscoacoustic Medium
- MS in mathematics from the University of Washington in 2003 Advisor: E. L. Stout
- BS in mathematics from New Mexico Tech in 2000 Advisor: D. R. Arterburn



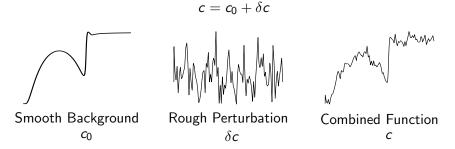




- The limitations of linearization
- The abstract forward model
 - Work based on the second-order equation of Lions-Magenes [Lions and Magenes, 1972, Non-homogeneous boundary value problems and applications, vol. 1]
- Continuous dependence on coefficients
 - Second-order equation proven by Stolk [Stolk, 2000, On the Modeling and Inversion of Seismic Data]
- The future: using the abstract forward model for general inversion, starting with 1-D



Most techniques used to solve inverse problems assume the functions describing the medium are oscillatory perturbations around a smooth medium.







Linearization depends on the separation of the medium into a low frequency (smooth) component and a high frequency (oscillatory) component. This does not match with reality, where there is no separation of scales.



Information at all frequencies





- The good news: there are projects currently underway in numerical nonlinear inversion (Dong Sun)
- The bad news: there is no theoretical framework justifying the ability to invert for nonsmooth media
- Even worse news: not even in one dimension
 - The *H*¹ theory and the work of Bube on discontinuous media is insufficient for general discontinuous media.





This project is the first step towards full nonlinear inversion.

- Analysis of the forward problem for general coefficients (L^{∞})
 - Existence proof = convergence of finite element method
 - Continuity of solutions w.r.t. coefficients
- The model is abstract enough that it covers many seismic models
 - Acoustics
 - Elastics
 - Viscoelastics





We consider the differential equation

$$Au' + Du + Bu + R[u] = f \in L^2(\mathbb{R}, H)$$

H is a Hilbert space (like L^2 , functions with finite energy)

- $A \in \mathcal{B}(H)$ is self-adjoint and positive-definite
- D skew-adjoint with dense domain $V \in H$
- $B \in \mathcal{B}(H)$
- $R[u](t) = \int Q(t-s)u(s) ds$, where $Q \in C(\mathbb{R}, \mathcal{B}(H))$ and Q(t) = 0 for t < 0



Consider the acoustic wave equation on a domain $\Omega \in \mathbb{R}^3$

$$\rho \frac{\partial \mathbf{v}}{\partial t} = -\nabla \boldsymbol{p} + \mathbf{f},$$

$$\frac{1}{\kappa} \frac{\partial \boldsymbol{p}}{\partial t} = -\nabla \cdot \mathbf{v}.$$

Define $H = (L^2(\Omega))^4$. Then $u = (p, v_1, v_2, v_3)^T$,

$$A = \begin{pmatrix} \frac{1}{\kappa} & 0 & 0 & 0\\ 0 & \rho & 0 & 0\\ 0 & 0 & \rho & 0\\ 0 & 0 & 0 & \rho \end{pmatrix}, \quad D = \begin{pmatrix} 0 & \frac{\partial}{\partial x_1} & \frac{\partial}{\partial x_2} & \frac{\partial}{\partial x_3}\\ \frac{\partial}{\partial x_1} & 0 & 0 & 0\\ \frac{\partial}{\partial x_2} & 0 & 0 & 0\\ \frac{\partial}{\partial x_3} & 0 & 0 & 0 \end{pmatrix}$$





Under assumptions of linearity and causality, the viscoelastic wave equation is

$$\rho \frac{\partial \mathbf{v}_i}{\partial t} = \sum_j \frac{\partial \sigma_{ij}}{\partial x_j} + f_i$$
$$\frac{\partial \sigma_{kl}}{\partial t} = \sum_{i,j} C_{ijkl} *_t \frac{1}{2} \left(\frac{\partial \mathbf{v}_i}{\partial x_j} + \frac{\partial \mathbf{v}_j}{\partial x_i} \right)$$

One choice of C is $C(\mathbf{x}, t) = \widetilde{C}(\mathbf{x})(\delta(t) - a(\mathbf{x})e^{-\alpha(\mathbf{x})t}H(t))$. δ is the Dirac delta, H is the Heaviside function, and $*_t$ denotes convolution in time.





Moving the convolution from the spatial derivatives to the time derivatives, we get

$$f_{i} = \rho(\mathbf{x}) \frac{\partial v_{i}}{\partial t} - \sum_{j} \frac{\partial \sigma_{ij}}{\partial x_{j}}$$

$$0 = \sum_{i,j} \widehat{C}_{ijkl}(\mathbf{x}) \frac{\partial \sigma_{ij}}{\partial t} - \frac{1}{2} \left(\frac{\partial v_{k}}{\partial x_{l}} + \frac{\partial v_{l}}{\partial x_{k}} \right)$$

$$+ \sum_{i,j} b_{ijkl}(\mathbf{x}) \sigma_{ij} + q_{ijkl}(\mathbf{x}, t) *_{t} \sigma_{ij}.$$

which fits with into the abstract model in a natural way

$$Au' + Du + Bu + R[u] = f$$





Our standard examples have $H = L^2$.

If a function is in L^{∞} , then the operator on L^2 given by multiplication against that function is a bounded operator.

$$\|fg\|_2 \le \|f\|_\infty \|g\|_2$$

So the abstract equation covers the case where we are trying to solve a differential equation with L^{∞} coefficients for an L^2 solution.





Theorem

A unique causal solution to the differential equation

Au' + Du + Bu + R[u] = f

exists provided that $f \in L^2(\mathbb{R}, H)$ is causal: $\operatorname{supp} f \subset [T_0, \infty)$ for some $T_0 \in \mathbb{R}$, and that the causal convolution kernel $Q \in L^1(\mathbb{R}, \mathcal{B}(H))$ is continuous in \mathbb{R}_+ : $Q \in C^0(\mathbb{R}_+, \mathcal{B}(H))$. The solution $u \in L^2(\mathbb{R}, H)$ and $\operatorname{supp} u \subset [T_0, \infty)$.





Let $\{w_k\}_{k=1}^{\infty} \subset V$ form a basis for H in V. Define the functions

$$u_m(t) = \sum_{k=1}^m g_{km}(t) w_k,$$

where the g_{km} 's are determined by the differential equation

$$\begin{array}{rcl} \langle u'_m(t),Aw_l\rangle - \langle u_m(t),Dw_l\rangle \\ + \langle u_m(t),B^*w_l\rangle + \langle u_m(t),R^*[w_l](t)\rangle & = & \langle f(t),w_l\rangle, & 1 \leq l \leq m, \\ u_m & = & 0 \text{ for } t \leq T_0. \end{array}$$





Convergence of the finite element approximations follow from energy estimates.

For the physical models, the abstract energy used here is the physical energy of the system.





If we have a sequence of equations

$$A_m u'_m + D_m u_m + B_m u_m + R_m [u_m] = f$$

and the coefficients converge in the weak sense

$$\lim_{m\to\infty} \|(A_m-A)w\|\to 0 \text{ for all } w\in H$$

Then u_m converges in measure.

If $H = L^2(\mathbb{R}^n)$ and the coefficients are L^{∞} , then L^1 convergence of the coefficients gives strong convergence of the solutions.





We can also take the derivative of the solution to the differential equation with respect to the coefficients.

If u_h is the solution to the differential equation with coefficients

$$A_h = A + h\delta A, B_h = B + h\delta B, Q_h = Q + h\delta Q$$

Then $(u_h - u)/h$ converges to the directional derivative of u in L^2 .





With existence, uniqueness, and convergence with respect to coefficients out of the way, the next step is to head towards nonlinear inversion of the 1-D problem.

What do we know so far?

• There is a one-to-one correspondence between H^1 impedances and L^2 impulse responses *h* which satisfy the acoustic transparency property

$$\langle f, h * f \rangle \geq \epsilon \|f\|^2$$

- Impedances which are functions of bounded variation satisfy acoustic transparency
- There exist non-BV functions which are not transparent





Based on what we know, we make the following conjecture:

- The natural realm of inversion for one-dimension is bounded variation
 - If the impedance is not BV, then transparency will fail
 - If the impedance is BV, then nonlinear inversion is possible

We hope to approach these problems using the convergence results for the abstract problem.





- We have shown that first-order integro-differential equations with coefficients forming bounded operators on Hilbert spaces have unique solutions in an appropriate sense.
- These equations include the acoustic wave equation, the elastic wave equation, and the viscoelastic wave equation with discontinuous coefficients as special cases.
- These solutions are continuous with respect to all parameters, that is, the coefficients of the equation, the initial condition, and the forcing function.
- We hope to use these results to establish nonlinear inversion for the one-dimensional problem





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