

Muhong Zhou

Rice University, Houston, TX

2011.8 - present

M.A., Ph.D. candidate in Computational and Applied Mathematics

- Ph.D.: *A Stable Composite Staggered Grid Finite Difference Scheme for Anisotropic Elastic Wave Simulations*
- M.A.: *Wave Equation Based Stencil Optimizations on a Multi-Core CPU*, 2014

Zhejiang University, Hangzhou, China

2007.8 - 2011.7

B.Sc. in Mathematics and Applied Mathematics

A Stable Composite Staggered Grid Finite Difference Scheme for Elastic Wave Simulations (TRIP Annual Review Meeting)

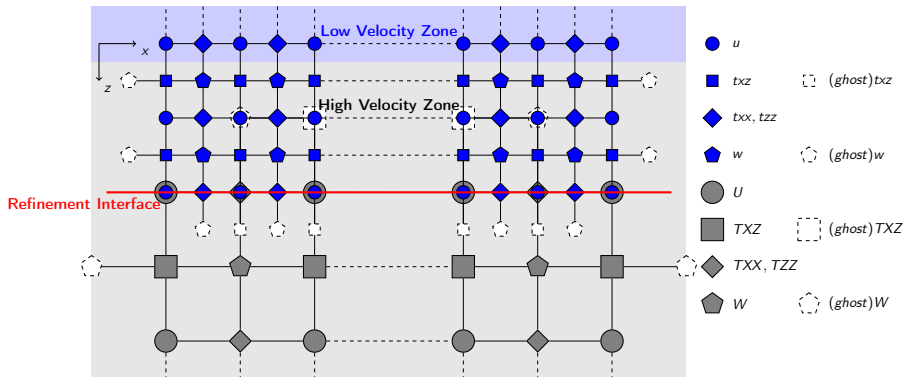
Muhong Zhou

CAAM, Rice University

April 25, 2016

I. Motivation of Using a Composite Grid

- Composite grid scheme uses less computing and memory resources than the uniform grid scheme
- Using a coarse grid in a high velocity region will not cause grid dispersion



II. Motivation of Seeking Stability

- Numerical stability: output error is bounded by the input error
- Lax-Richtmyer Equivalence Theorem: For a linear consistent scheme, stability and convergence are equivalent
- Tool to build stability: [energy method](#)

Previous Work and My Contribution

Previous papers on the energy method to build stability:

- [Pettersson and Sjögreen, 2010]¹ finite difference scheme, composite collocated grid, second order wave equation
- [Rodríguez, 2008]² finite element scheme, composite staggered grid, first order wave equation system

My contribution:

- use the energy method to build a stable composite staggered grid finite difference scheme for first order wave equation system

¹Pettersson and Sjögreen, *Stable grid refinement and singular source discretization for seismic wave simulations*, 2010

²Rodríguez, *A spurious-free space-time mesh refinement for elastodynamics*, 2008

Previous Work and My Contribution

Previous papers on the energy method to build stability:

- [Pettersson and Sjögreen, 2010]¹ **finite difference scheme**, composite collocated grid, second order wave equation
- [Rodríguez, 2008]² finite element scheme, **composite staggered grid**, **first order wave equation system**

My contribution:

- use the energy method to build a stable composite staggered grid finite difference scheme for first order wave equation system

¹Pettersson and Sjögreen, *Stable grid refinement and singular source discretization for seismic wave simulations*, 2010

²Rodríguez, *A spurious-free space-time mesh refinement for elastodynamics*, 2008

Steps to Use the Energy Method to Get Stability

- 1 Formulate the energy: $E[s_1, \dots, s_m](t)$,
where $s_i(t) (i = 1, \dots, m)$ denote simulation variables

When the source is absent, make sure:

- 2 the energy is non increasing over time:

$$E[s_1, \dots, s_m](t_2) \leq E[s_1, \dots, s_m](t_1), \text{ for } t_2 \geq t_1$$

- 3 the energy is bounded above and below by simulation variables:

$$a \sum_{i=1}^m \|s_i(t)\|^2 \leq E[s_1, \dots, s_m](t) \leq b \sum_{i=1}^m \|s_i(t)\|^2, \quad a, b > 0$$

-
- 2 and 3 yield stability:

$$a \sum_{i=1}^m \|s_i(T)\|^2 \stackrel{3}{\leq} E[s_1, \dots, s_m](T) \stackrel{2}{\leq} E[s_1, \dots, s_m](0) \stackrel{3}{\leq} b \sum_{i=1}^m \|s_i(0)\|^2$$

2D Example

Isotropic elastic wave equation in first order form:

$$\rho \dot{u} = t_{xx} + t_{xz} + s_1$$

$$\rho \dot{w} = t_{zx} + t_{zz} + s_2$$

$$\dot{t}_{xx} = (\lambda + 2\mu)u_x + \lambda w_z$$

$$\dot{t}_{zz} = \lambda u_x + (\lambda + 2\mu)w_z$$

$$\dot{t}_{xz} = \mu(u_z + w_x)$$

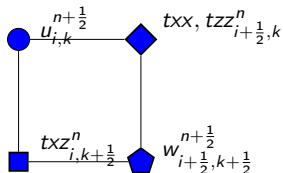
Energy in the continuous time-space:

$$\text{Kinetic energy: } k(t) = \frac{1}{2} \iint \rho(u^2 + w^2) dx dz$$

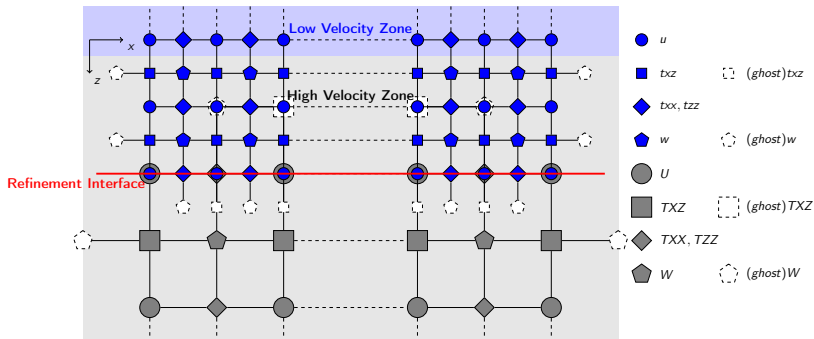
$$\text{Strain energy: } s(t) = \frac{1}{2} \iint t_{xx} \cdot \epsilon_{xx} + t_{zz} \cdot \epsilon_{zz} + 2t_{xz} \cdot \epsilon_{xz} dx dz$$

When the source is absent, $e(t) = k(t) + s(t)$ satisfies $\dot{e}(t) = 0$.

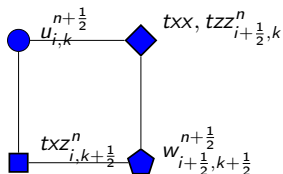
A 2D Composite Staggered Grid with (2,2) FDM



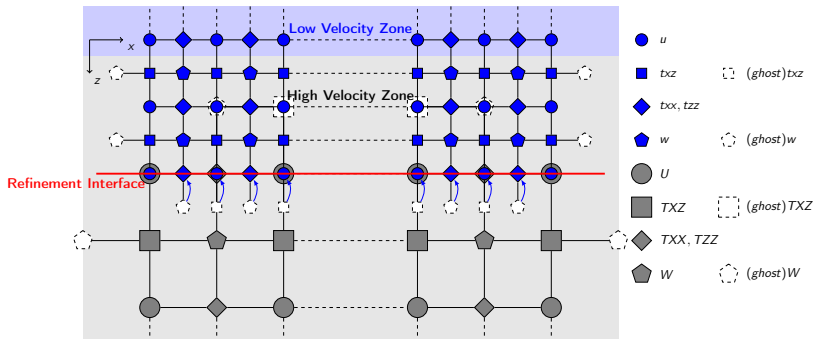
$$\left\{ \begin{array}{ll} u_{i,k}^{n+\frac{1}{2}} & i = 0, \dots, nx; k = 0, \dots, nz \\ t_{xx}, t_{zz}^n_{i+\frac{1}{2},k} & i = 0, \dots, nx - 1; k = 0, \dots, nz \\ t_{xz}^n_{i,k+\frac{1}{2}} & i = 0, \dots, nx; k = 0, \dots, nz - 1 \\ w_{i+\frac{1}{2},k+\frac{1}{2}}^{n+\frac{1}{2}} & i = 0, \dots, nx - 1; k = 0, \dots, nz - 1 \end{array} \right.$$



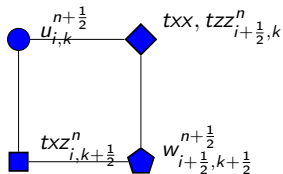
A 2D Composite Staggered Grid with (2,2) FDM



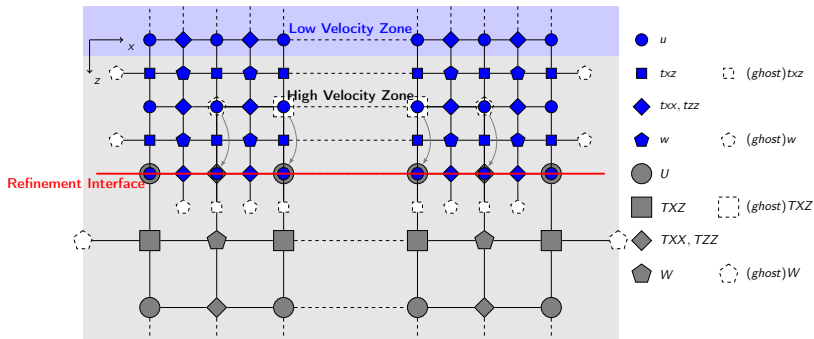
$$\left\{ \begin{array}{ll} u_{i,k}^{n+\frac{1}{2}} & i = 0, \dots, nx; k = 0, \dots, nz \\ t_{xx}, t_{zz}_{i+\frac{1}{2},k}^n & i = 0, \dots, nx - 1; k = 0, \dots, nz \\ t_{xz}_{i,k+\frac{1}{2}}^n & i = 0, \dots, nx; k = 0, \dots, nz - 1 \\ w_{i+\frac{1}{2},k+\frac{1}{2}}^{n+\frac{1}{2}} & i = 0, \dots, nx - 1; k = 0, \dots, nz - 1 \end{array} \right.$$



A 2D Composite Staggered Grid with (2,2) FDM



$$\left\{ \begin{array}{ll} u_{i,k}^{n+\frac{1}{2}} & i = 0, \dots, nx; k = 0, \dots, nz \\ t_{xx}, t_{zz}^n_{i+\frac{1}{2},k} & i = 0, \dots, nx - 1; k = 0, \dots, nz \\ t_{xz}^n_{i,k+\frac{1}{2}} & i = 0, \dots, nx; k = 0, \dots, nz - 1 \\ w_{i+\frac{1}{2},k+\frac{1}{2}}^{n+\frac{1}{2}} & i = 0, \dots, nx - 1; k = 0, \dots, nz - 1 \end{array} \right.$$



Define Discrete Energy $e^n = k^n + s^n$

Define inner products (or quadrature rules) on the grid:

$$\langle u^1, u^2 \rangle_u = h^2 \sum_{i=0}^{nx} \sum_{k=0}^{nz} \alpha_i \alpha_k u_{i,k}^1 u_{i,k}^2 \approx \iint u^1 u^2 dx dz$$

$$\langle w^1, w^2 \rangle_w = h^2 \sum_{i=0}^{nx-1} \sum_{k=0}^{nz-1} w_{i+\frac{1}{2}, k+\frac{1}{2}}^1 w_{i+\frac{1}{2}, k+\frac{1}{2}}^2 \approx \iint w^1 w^2 dx dz$$

$$\langle txz^1, txz^2 \rangle_{txz} = h^2 \sum_{i=0}^{nx} \sum_{k=0}^{nz-1} \alpha_i txz_{i, k+\frac{1}{2}}^1 txz_{i, k+\frac{1}{2}}^2 \approx \iint txz^1 txz^2 dx dz$$

where $\alpha_{i/k} = \frac{1}{2}$ at two ends of axis i/k , and 1 elsewhere.

$$k^n = \frac{1}{2} \langle \rho u^{n+\frac{1}{2}}, u^{n-\frac{1}{2}} \rangle_u + \frac{1}{2} \langle \rho w^{n+\frac{1}{2}}, w^{n-\frac{1}{2}} \rangle_w \approx \frac{1}{2} \iint \rho (u^2 + w^2) dx dz$$

$$s^n = \frac{1}{2} h^2 \sum_{i=0}^{nx-1} \sum_{k=0}^{nz} \alpha_k \begin{pmatrix} txx_{i+\frac{1}{2}, k}^n \\ tzz_{i+\frac{1}{2}, k}^n \end{pmatrix}^T \begin{pmatrix} (\lambda + 2\mu)_{i+\frac{1}{2}, k} & \lambda_{i+\frac{1}{2}, k} \\ \lambda_{i+\frac{1}{2}, k} & (\lambda + 2\mu)_{i+\frac{1}{2}, k} \end{pmatrix}^{-1} \begin{pmatrix} txx_{i+\frac{1}{2}, k}^n \\ tzz_{i+\frac{1}{2}, k}^n \end{pmatrix} \\ + \frac{1}{2} \langle \frac{1}{\mu} txz^n, txz^n \rangle_{txz} \approx \frac{1}{2} \iint txx \cdot \epsilon_{xx} + tzz \cdot \epsilon_{zz} + 2txz \cdot \epsilon_{xz} dx dz$$

Define Discrete Energy $e^n = k^n + s^n$

Define inner products (or quadrature rules) on the grid:

$$\langle u^1, u^2 \rangle_u = h^2 \sum_{i=0}^{nx} \sum_{k=0}^{nz} \alpha_i \alpha_k u_{i,k}^1 u_{i,k}^2 \approx \iint u^1 u^2 dx dz$$

$$\langle w^1, w^2 \rangle_w = h^2 \sum_{i=0}^{nx-1} \sum_{k=0}^{nz-1} w_{i+\frac{1}{2}, k+\frac{1}{2}}^1 w_{i+\frac{1}{2}, k+\frac{1}{2}}^2 \approx \iint w^1 w^2 dx dz$$

$$\langle txz^1, txz^2 \rangle_{txz} = h^2 \sum_{i=0}^{nx} \sum_{k=0}^{nz-1} \alpha_i txz_{i, k+\frac{1}{2}}^1 txz_{i, k+\frac{1}{2}}^2 \approx \iint txz^1 txz^2 dx dz$$

where $\alpha_{i/k} = \frac{1}{2}$ at two ends of axis i/k , and 1 elsewhere.

$$k^n = \frac{1}{2} \langle \rho u^{n+\frac{1}{2}}, u^{n-\frac{1}{2}} \rangle_u + \frac{1}{2} \langle \rho w^{n+\frac{1}{2}}, w^{n-\frac{1}{2}} \rangle_w \approx \frac{1}{2} \iint \rho (u^2 + w^2) dx dz$$

$$s^n = \frac{1}{2} h^2 \sum_{i=0}^{nx-1} \sum_{k=0}^{nz} \alpha_k \begin{pmatrix} txx_{i+\frac{1}{2}, k}^n \\ tzz_{i+\frac{1}{2}, k}^n \end{pmatrix}^T \begin{pmatrix} (\lambda + 2\mu)_{i+\frac{1}{2}, k} & \lambda_{i+\frac{1}{2}, k} \\ \lambda_{i+\frac{1}{2}, k} & (\lambda + 2\mu)_{i+\frac{1}{2}, k} \end{pmatrix}^{-1} \begin{pmatrix} txx_{i+\frac{1}{2}, k}^n \\ tzz_{i+\frac{1}{2}, k}^n \end{pmatrix} \\ + \frac{1}{2} \langle \frac{1}{\mu} txz^n, txz^n \rangle_{txz} \approx \frac{1}{2} \iint txx \cdot \epsilon_{xx} + tzz \cdot \epsilon_{zz} + 2txz \cdot \epsilon_{xz} dx dz$$

Energy Conservation Condition

1 Compute energy difference on the fine grid $e^n - e^{n-1}$:

$$h\Delta t \sum_{i=1}^{nx-1} u_{i,nz}^{n-\frac{1}{2}} \frac{\overline{txz_{i,nz}^n} + \overline{txz_{i,nz}^{n-1}}}{2} + h\Delta t \sum_{i=0}^{nx-1} w_{i+\frac{1}{2},nz}^{n-\frac{1}{2}} \frac{tzz_{i+\frac{1}{2},nz}^n + tzz_{i+\frac{1}{2},nz}^{n-1}}{2}$$

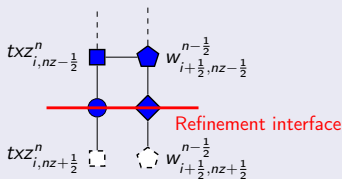
2 Compute energy difference on the coarse grid $E^n - E^{n-1}$:

$$-H\Delta t \sum_{i=1}^{NX-1} U_{i,0}^{n-\frac{1}{2}} \frac{\overline{TXZ_{i,0}^n} + \overline{TXZ_{i,0}^{n-1}}}{2} - H\Delta t \sum_{i=0}^{nx-1} W_{i+\frac{1}{2},0}^{n-\frac{1}{2}} \frac{TZZ_{i+\frac{1}{2},0}^n + TZZ_{i+\frac{1}{2},0}^{n-1}}{2}$$

Averaging operator ($\bar{\cdot}$):

$$\overline{txz_{i,nz}^n} = \frac{1}{2} (txz_{i,nz+\frac{1}{2}}^n + txz_{i,nz-\frac{1}{2}}^n)$$

$$w_{i+\frac{1}{2},nz}^{n-\frac{1}{2}} = \frac{1}{2} (w_{i+\frac{1}{2},nz+\frac{1}{2}}^{n-\frac{1}{2}} + w_{i+\frac{1}{2},nz-\frac{1}{2}}^{n-\frac{1}{2}})$$



Energy Conservation Condition

1 Compute energy difference on the fine grid $e^n - e^{n-1}$:

$$h\Delta t \sum_{i=1}^{nx-1} u_{i,nz}^{n-\frac{1}{2}} \frac{\overline{txz_{i,nz}^n} + \overline{txz_{i,nz}^{n-1}}}{2} + h\Delta t \sum_{i=0}^{nx-1} w_{i+\frac{1}{2},nz}^{n-\frac{1}{2}} \frac{tzz_{i+\frac{1}{2},nz}^n + tzz_{i+\frac{1}{2},nz}^{n-1}}{2}$$

2 Compute energy difference on the coarse grid $E^n - E^{n-1}$:

$$-H\Delta t \sum_{i=1}^{NX-1} U_{i,0}^{n-\frac{1}{2}} \frac{\overline{TXZ_{i,0}^n} + \overline{TXZ_{i,0}^{n-1}}}{2} - H\Delta t \sum_{i=0}^{nx-1} W_{i+\frac{1}{2},0}^{n-\frac{1}{2}} \frac{TZZ_{i+\frac{1}{2},0}^n + TZZ_{i+\frac{1}{2},0}^{n-1}}{2}$$

Energy Conservation Condition

Then the total energy is conserved (i.e., $e^n + E^n = e^{n-1} + E^{n-1}$) as long as for any m, n ,

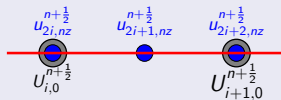
$$h \sum_{i=1}^{nx-1} u_{i,nz}^{n-\frac{1}{2}} \overline{txz_{i,nz}^m} = H \sum_{i=1}^{NX-1} U_{i,0}^{n-\frac{1}{2}} \overline{TXZ_{i,0}^m}$$
$$h \sum_{i=0}^{nx-1} w_{i+\frac{1}{2},nz}^{n-\frac{1}{2}} tzz_{i+\frac{1}{2},nz}^m = H \sum_{i=0}^{NX-1} W_{i+\frac{1}{2},0}^{n-\frac{1}{2}} TZZ_{i+\frac{1}{2},0}^m$$

Energy Conserving Interpolation

Assume $H=2h$, decouple
$$h \sum_{i=1}^{nx-1} u_{i,nz}^{n-\frac{1}{2}} \overline{txz_{i,nz}^m} = H \sum_{i=1}^{NX-1} U_{i,0}^{n-\frac{1}{2}} \overline{TXZ_{i,0}^m} \quad (*)$$

Transmission condition (u, w, txz, tzz):

$$1 \begin{cases} u_{2i,nz}^{n+\frac{1}{2}} = U_{i,0}^{n+\frac{1}{2}} \\ u_{2i+1,nz}^{n+\frac{1}{2}} = \frac{1}{2}(U_{i,0}^{n+\frac{1}{2}} + U_{i+1,0}^{n+\frac{1}{2}}) \end{cases}$$



Plugging 1 back into (*) yields

$$2 \overline{TXZ_{i,0}^n} = \overline{txz_{2i,nz}^n} + \frac{1}{2}(\overline{txz_{2i+1,nz}^n} + \overline{txz_{2i-1,nz}^n})$$

$$\begin{array}{c} \boxed{} TXZ_{i,-\frac{1}{2}}^{n+\frac{1}{2}} \\ \\ \begin{array}{ccc} \blacksquare & \blacksquare & \blacksquare \\ \overline{txz_{2i-1,nz-\frac{1}{2}}^n} & \overline{txz_{2i,nz-\frac{1}{2}}^n} & \overline{txz_{2i+1,nz-\frac{1}{2}}^n} \\ \hline \overline{txz_{2i-1,nz+\frac{1}{2}}^n} & \overline{txz_{2i,nz+\frac{1}{2}}^n} & \overline{txz_{2i+1,nz+\frac{1}{2}}^n} \\ \boxed{} & \boxed{} & \boxed{} \end{array} \\ \\ \blacksquare TXZ_{i,\frac{1}{2}}^n \end{array}$$

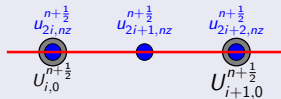
(Energy conserving interpolation): 1 and 2 together form a strictly diagonally dominant linear equation system for ghost data ($txz_{i,nz+\frac{1}{2}}^n, TXZ_{i,-\frac{1}{2}}^n$).

Energy Conserving Interpolation

Assume $H=2h$, decouple
$$h \sum_{i=1}^{nx-1} u_{i,nz}^{n-\frac{1}{2}} \overline{txz_{i,nz}^m} = H \sum_{i=1}^{NX-1} U_{i,0}^{n-\frac{1}{2}} \overline{TXZ_{i,0}^m} \quad (*)$$

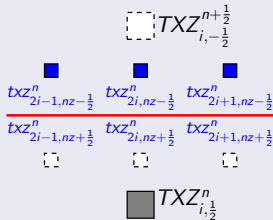
Transmission condition (u, w, txz, tzz):

$$1 \begin{cases} u_{2i,nz}^{n+\frac{1}{2}} = U_{i,0}^{n+\frac{1}{2}} \\ u_{2i+1,nz}^{n+\frac{1}{2}} = \frac{1}{2}(U_{i,0}^{n+\frac{1}{2}} + U_{i+1,0}^{n+\frac{1}{2}}) \end{cases}$$



Plugging 1 back into (*) yields

$$2 \overline{TXZ_{i,0}^n} = \overline{txz_{2i,nz}^n} + \frac{1}{2}(\overline{txz_{2i+1,nz}^n} + \overline{txz_{2i-1,nz}^n})$$



(**Energy conserving interpolation**): 1 and 2 together form a strictly diagonally dominant linear equation system for ghost data ($txz_{i,nz+\frac{1}{2}}^n, TXZ_{i,-\frac{1}{2}}^n$).

Qualify for Boundedness

With energy conserving interpolation, Cauchy-Schwarz inequality

$$E^n + e^n \geq \underbrace{\frac{1}{2}(\langle \rho u^{n-\frac{1}{2}}, u^{n-\frac{1}{2}} \rangle_u - \frac{\Delta t(\gamma_1 + \gamma_2)}{h} \|u^{n-\frac{1}{2}}\|_u^2)}_{\text{implies } \Delta t < \frac{h\rho_{\min}}{\gamma_1 + \gamma_2}} + \underbrace{\frac{1}{2}(\langle \frac{t_{xx}^n}{2(\lambda + \mu)}, t_{xx}^n \rangle_{t_{xx}} - \frac{\Delta t}{h\gamma_1} \|t_{xx}^n\|_{t_{xx}}^2)}_{\text{implies } \Delta t < \frac{h\gamma_1}{2(\lambda + \mu)_{\max}}} + \dots$$

$\gamma_1, \gamma_2, \dots$: free parameters resulted from C-S, have same unit as $\sqrt{\kappa\rho}$

$$\Delta t < \max_{\gamma_1, \gamma_2, \dots} \left\{ \min \left\{ \frac{h\rho_{\min}}{\gamma_1 + \gamma_2}, \frac{h\gamma_1}{2(\lambda + \mu)_{\max}}, \dots \right\} \right\}$$

Qualify for Boundedness

$$\Delta t < \max_{\gamma_1, \gamma_2, \dots} \left\{ \min \left\{ \frac{h\rho_{\min}}{\gamma_1 + \gamma_2}, \frac{h\gamma_1}{2(\lambda + \mu)_{\max}}, \dots \right\} \right\}$$

- Sufficient condition for stability vs. CFL condition (necessary condition for stability)
- Depends on both density and velocity vs. CFL condition (depends on velocity only)
- In isotropic homogenous medium, this upper bound reduces to CFL condition:

$$\Delta t < \max_{\gamma_1, \gamma_2} \left\{ \min \left\{ \frac{h\rho}{\gamma_1 + \gamma_2}, \frac{h\gamma_1}{2(\lambda + \mu)}, \frac{h\gamma_2}{2\mu} \right\} \right\} = h \sqrt{\frac{\rho}{2(\lambda + 2\mu)}} = \frac{h}{\sqrt{2} V_p} \quad (\text{CFL})$$

Results: Compare Horizontal Velocity

- 20 Hz Ricker wavelet; delay time: 0.05 s
- $V_p = 1500$ m/s, $V_s = 1000$ m/s, $\rho = 1000$ kg/m³
- Snapshot is taken at $T = 0.2$ s

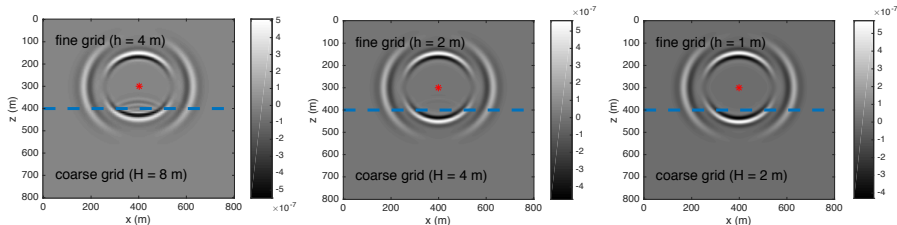
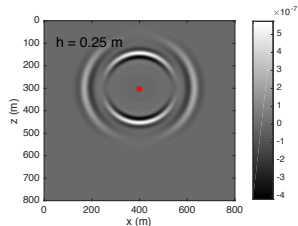
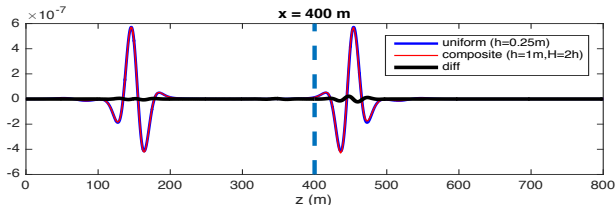
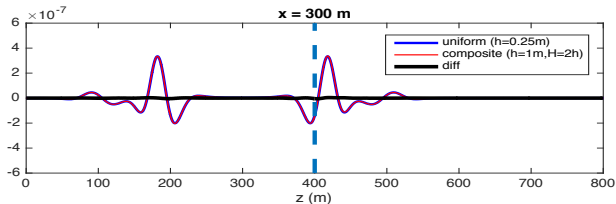
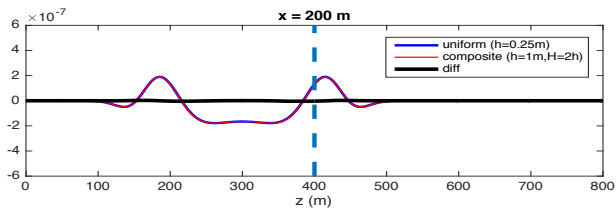
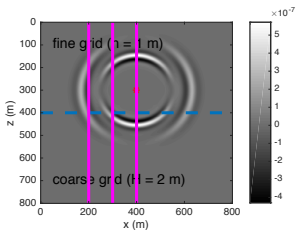


Figure: [top right] uniform grid solution ($h=0.25$ m);
[bottom] composite grid solutions ($h = 4, 2, 1$ m, $H=2h$).

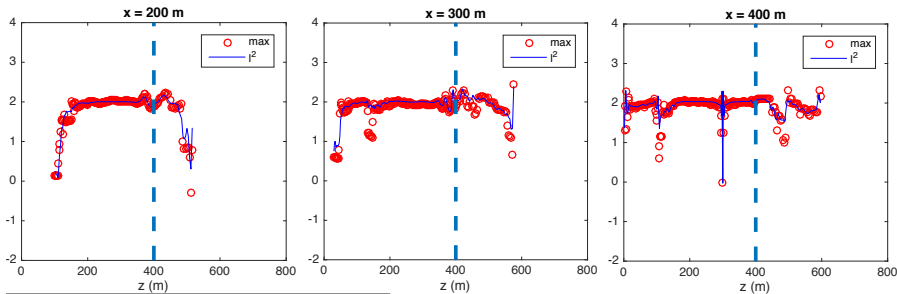
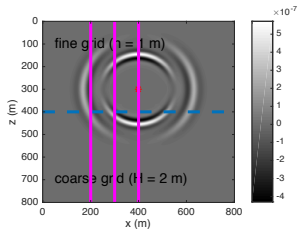
Results: Compare Vertical Traces



Results: Estimate Convergence Rate³ (Tails Truncated)

- Estimate convergence rate at $x = 200, 300, 400\text{m}$
- $R(z) = \log_2\left(\frac{\|u_h(z) - u_{h/2}(z)\|}{\|u_{h/2}(z) - u_{h/4}(z)\|}\right)$ with $h = 2, 1, 0.5\text{ m}$

where $\|u(z)\|_{l^2} = \sqrt{\sum_k |u(z, t_k)|^2}$ or
 $\|u(z)\|_{\infty} = \max_k |u(z, t_k)|$.

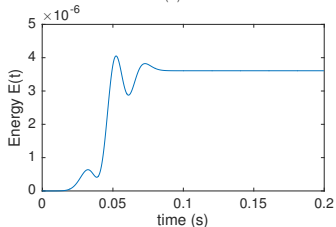
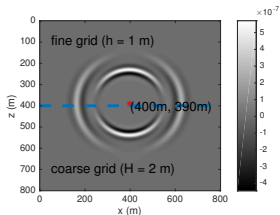


³Bencomo, *Joint model and minimal-source full waveform inversion for seismic imaging under general anisotropic sources*, 2015

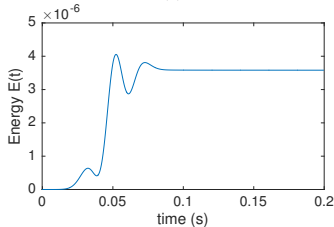
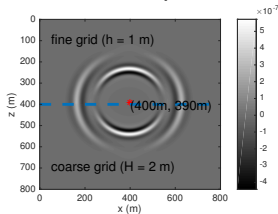
Results: Energy Conserving vs. Intuitive Interpolation

Source: (400 m, 390 m)

Energy Conserving Interpolation



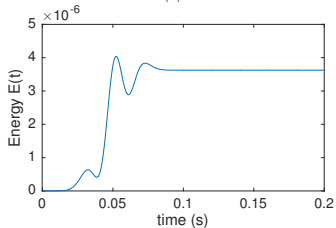
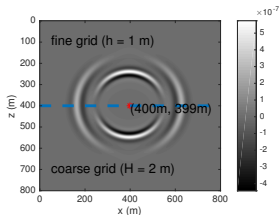
Intuitive Interpolation



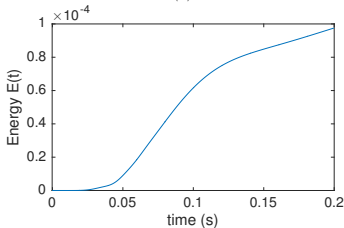
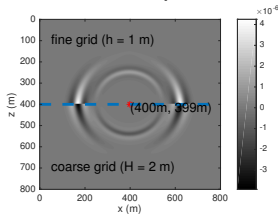
Results: Energy Conserving vs. Intuitive Interpolation

Source: (400 m, 399 m)

Energy Conserving Interpolation



Intuitive Interpolation



Summary and Future Work

Summary

- Proposed a numerically stable composite staggered grid FDM
- Achieved stability by implementing the energy conserving interpolation on the interface ghost data
- Demonstrated an upper bound on the time step
- Numerically verified the convergence rate for (2,2) scheme

Future Work

- Load into IWAVE, parallelization, performance tuning
- Elasto-acoustic coupling
- Towards high-order scheme

Thank You!
Q&A